

考虑桩体固结和土体渗透性抛物线分布固结解

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摘 要: 在传统的复合地基固结理论中, 等应变条件和桩土流量连续的假设之间存在着不一致的地方。为此, 抛弃了桩土界面上水流量相等的假设, 考虑了桩体径向渗流引起的桩体固结和变形, 同时考虑了桩体施工扰动而造成的扰动区土体水平渗透系数抛物线分布的特点, 给出了此类固结问题的控制方程和解答以及复合地基固结度的表达式, 并对复合地基固结度的解进行了讨论, 可以看出无论井径比(影响区半径和桩体半径之比)趋于 1 或无穷大, 本文固结解均能退化到 Terzaghi 天然地基一维固结的解答。最后对本文的解和以往的解做了比较, 结果表明: 本文解给出的地基固结度比考虑扰动区土体水平渗透系数为直线变化或者保持不变的解给出的固结度更大; 随着井径比的减小, 本文解和以往的解得到的固结度差值变大。

关键词: 固结; 复合地基; 桩体固结; 水平渗透系数; 抛物线分布

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Solutions for consolidation considering column consolidation and parabolic distribution of soil permeability

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Abstract: In the traditional consolidation theories for composite foundation, there is an inconsistency between the equal strain assumption and the assumption of flow continuity at the column-soil interface. To remedy this inconsistency, the radial flow within the column instead of the equal flow assumption at the column-soil interface is considered to account for the consolidation and deformation of the column. In addition, the parabolic distribution of the horizontal permeability of soil in the disturbed zone due to the influence of the column construction is considered too. The governing equations for this type of consolidation problem are derived and the solutions to the governing equations are obtained. Then the solution of the average degree of consolidation is obtained and discussed. It can be seen that whenever the radius ratio (radius ratio of the influence zone to the column) approaches 1 or ∞ , the present solution for composite foundation all can be reduced to the solution of Terzaghi's one-dimensional consolidation for natural foundation. Finally, a comparison between the present solution and some previous solutions is made. The results show that the rate of consolidation calculated by the present solution is rapider than that by the solutions which assume the horizontal permeability of soil in the disturbed zone is linearly changed or kept constant as a reduced value. The difference between the average degree calculated by the present solution and that by the previous solutions increases with the reduction in the radius ratio.

Key words: consolidation; composite foundation; column consolidation; horizontal permeability; parabolic distribution

0 前 言

在以往的复合地基固结理论中, 均采用了桩土界面流量相等的假设, 即假定任一时刻任一深度处从土

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体流入桩体的水量等于桩体中向上水流的增量,也就是说桩体中的水量是假定保持不变的,进而也可以说桩体的体积被假定为是一个不变的量。然而,等应变条件假设桩体和土体的变形在任一深度处是一样。显然,这两个假定是相互矛盾的。鉴于此,本文抛弃了桩土界面流量相等的假定,通过考虑了桩体内的径向和竖向渗流,对桩体和土体均采用固结方程求解,认为桩体排出地基的水量不等于土体进入桩体的水量,而是等于土体进入桩体的水量和桩体体积变形量之和。这样,即可说明桩体内水量变化和桩体体积变化的关系,从而避免了桩土界面流量相等的假设和等应变假设的矛盾。另外,桩体的施工会对周围的土体产生扰动作用,传统的固结解^[1-3]在考虑扰动作用时均假定扰动区土体的渗透系数是一个小于未扰动区土体渗透系数的常量。然而,SHARMA 等^[4]的研究发现,由于桩体周边的土体受到的扰动程度不同,土体的渗透系数在扰动区是一个连续变化的量,而不是一个常量。针对这一点,ZHANG 等^[5-6]以及 Xie 等^[7]给出了一种考虑土体水平渗透系数呈直线变化的复合地基固结解,ROHAN^[8]则给出了一种考虑涂抹区水平渗透系数呈抛物线变化的砂井地基固结解,然而该文只限于考虑土体的径向渗流。Xie 等^[9]还对传统的桩周流量连续假定做了改进并得到了此类固结问题的控制方程及解答。

本文通过考虑桩体内的径竖向组合渗流,对桩体采用固结方程求解,同时考虑了桩体施工扰动造成的扰动区土体水平渗透系数抛物线变化的特点,给出了此类固结问题的控制方程和解答,并给出了复合地基固结度的表达式,最后对本文解和以往解做了比较。

1 定解问题

1.1 控制方程及求解条件的推导

图 1 为本文采用的固结简化模型,可以看出,和以往的模型不同的是,本文考虑了桩体内的径向渗流,这样桩体内的孔压在任一深度处将不是一个常量,而是随径向距离变化的量。

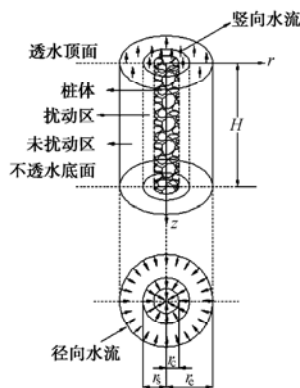


图 1 固结简化模型

Fig. 1 Simplified model for consolidation

本文在推导时做了如下的假定:

- (1) 等应变条件成立。
- (2) 土中水的渗流服从 Darcy 定律。
- (3) 外部荷载一次瞬时施加,荷载在复合地基中引起的附加应力沿深度均匀分布,即 $\sigma = \sigma_0$ 。
- (4) 假定扰动区土体的水平渗透系数随着径向距离按抛物线变化,如图 2 所示。

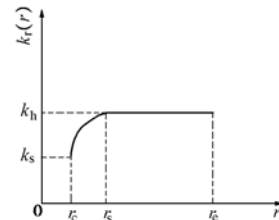


图 2 土体水平渗透系数变化曲线

Fig. 2 Variation of horizontal permeability of soil

参考 Rohan^[8]的研究,扰动区土体水平渗透系数 $k_r(r)$ 必须满足以下 3 个条件:

$$\left. \begin{aligned} k_r(r_c) &= k_s, \\ k_r(r_s) &= k_h, \\ \frac{dk_r(r_s)}{dr} &= 0, \end{aligned} \right\} \quad (1)$$

满足上述条件的扰动区土体的水平渗透系数可以表示为

$$k_r(r) = k_h \left(1 - \frac{1}{\alpha} \right) \left(A - B + C \frac{r}{r_c} \right) \left(A + B - C \frac{r}{r_c} \right). \quad (2)$$

式中 $\alpha = \frac{k_h}{k_s}$, $A = \sqrt{\frac{\alpha}{\alpha-1}}$, $B = \frac{s}{s-1}$, $C = \frac{1}{s-1}$; r_c , r_s , r_e 分别为桩体、扰动区和整个影响区的半径; $k_r(r)$ 和 k_s 分别为扰动区任一点和桩土界面处的渗透系数,如图 2 所示; k_h 为未扰动区土体的水平渗透系数; $s = r_s/r_c$ 。

再由平衡方程和等应变条件有:

$$\pi(r_e^2 - r_c^2) \bar{\sigma}_s + \pi r_c^2 \bar{\sigma}_c = \pi r_e^2 \sigma_0, \quad (3)$$

$$\frac{\bar{\sigma}_s - \bar{u}_s}{E_s} = \frac{\bar{\sigma}_c - \bar{u}_c}{E_c} = \varepsilon_v. \quad (4)$$

式中 $\bar{\sigma}_s$, $\bar{\sigma}_c$ 分别为土体和桩体中的平均总应力; E_s , E_c 分别为土体和桩体的压缩模量; ε_v 为地基任一深度处的竖向应变; σ_0 为上部荷载在复合地基任一深度处引起的平均总附加应力,其值是一个常量。 $\bar{u}_s$, $\bar{u}_c$ 分别为土体和桩体内的任一深度处的平均超静孔压,其表达式为

$$\bar{u}_s = \frac{1}{\pi(r_e^2 - r_c^2)} \left(\int_{r_c}^{r_s} 2\pi r u_d dr + \int_{r_s}^{r_e} 2\pi r u_n dr \right), \quad (5)$$

$$\bar{u}_c = \frac{1}{\pi r_c^2} \int_0^{r_c} 2\pi r u_c dr, \quad (6)$$

式中, u_d , u_n , u_c 分别为扰动区和未扰动区土体以及桩体内任一点的超静孔压。

则由式(3)、(4)可得:

$$\varepsilon_v = \frac{n^2 \sigma_0 - [(n^2 - 1)\bar{u}_s + \bar{u}_c]}{E_s(n^2 - 1 + Y)} = \frac{n^2(\sigma_0 - \bar{u})}{E_s(n^2 - 1 + Y)}, \quad (7)$$

式中, $Y = \frac{E_c}{E_s}$, $n = \frac{r_c}{r_s}$, $\bar{u}$ 为地基任一深度处总的平均超静孔压, 其表达式可以根据面积加权的方法得到:

$$\begin{aligned} \bar{u} &= \frac{1}{\pi r_c^2} \left(\int_0^{r_c} 2\pi r u_c dr + \int_{r_c}^{r_s} 2\pi r u_d dr + \int_{r_s}^{r_c} 2\pi r u_n dr \right) \\ &= \frac{(n^2 - 1)\bar{u}_s + \bar{u}_c}{n^2}. \end{aligned} \quad (8)$$

则对(7)关于时间求导可得:

$$\frac{\partial \varepsilon_v}{\partial t} = -\frac{n^2}{E_s(n^2 - 1 + Y)} \frac{\partial \bar{u}}{\partial t}. \quad (9)$$

由 WANG 等^[31]和 ROHAN^[8]的研究, 桩体和土体的固结方程可以写为

$$\frac{1}{r} \frac{\partial}{\partial r} \left[\frac{k_{hc}}{\gamma_w} r \frac{\partial u_c}{\partial r} \right] + \frac{k_{vc}}{\gamma_w} \frac{\partial^2 \bar{u}_c}{\partial z^2} = -\frac{\partial \varepsilon_v}{\partial t} \quad (0 \leq r \leq r_c), \quad (10a)$$

$$\frac{1}{r} \frac{\partial}{\partial r} \left[\frac{k_r(r)}{\gamma_w} r \frac{\partial u_d}{\partial r} \right] + \frac{k_v}{\gamma_w} \frac{\partial^2 \bar{u}_s}{\partial z^2} = -\frac{\partial \varepsilon_v}{\partial t} \quad (r_c \leq r \leq r_s), \quad (10b)$$

$$\frac{1}{r} \frac{\partial}{\partial r} \left(\frac{k_h}{\gamma_w} r \frac{\partial u_n}{\partial r} \right) + \frac{k_v}{\gamma_w} \frac{\partial^2 \bar{u}_s}{\partial z^2} = -\frac{\partial \varepsilon_v}{\partial t} \quad (r_s \leq r \leq r_c), \quad (10c)$$

式中, k_v 分别为土体的竖向渗透系数, k_{hc} , k_{vc} 分别为桩体水平和竖向的渗透系数, γ_w 为水的重度。

径向求解条件为

$$\frac{\partial u_n}{\partial r} = 0 \quad (r = r_c), \quad (11a)$$

$$k_r(r_s) \frac{\partial u_d}{\partial r} = k_h \frac{\partial u_n}{\partial r} \quad (r = r_s), \quad (11b)$$

$$u_d = u_n \quad (r = r_s), \quad (11c)$$

$$u_d = u_c \quad (r = r_c), \quad (11d)$$

$$k_r(r_c) \frac{\partial u_d}{\partial r} = k_{hc} \frac{\partial u_c}{\partial r} \quad (r = r_c), \quad (11e)$$

$$\frac{\partial u_c}{\partial r} = 0 \quad (r = 0), \quad (11f)$$

对方程(10b)和(10c)两边关于 r 积分, 并应用边界条件(11a)和(11b)有:

$$\frac{\partial u_d}{\partial r} = \frac{\gamma_w}{2k_r(r)} \left(\frac{\partial \varepsilon_v}{\partial t} + \frac{k_v}{\gamma_w} \frac{\partial^2 \bar{u}_s}{\partial z^2} \right) \left(\frac{r_c^2}{r} - r \right), \quad (12a)$$

$$\frac{\partial u_n}{\partial r} = \frac{\gamma_w}{2k_h} \left(\frac{\partial \varepsilon_v}{\partial t} + \frac{k_v}{\gamma_w} \frac{\partial^2 \bar{u}_s}{\partial z^2} \right) \left(\frac{r_c^2}{r} - r \right). \quad (12b)$$

再对式(12a)、(12b)两边关于 r 积分并应用边界条件(11c)、(11d)有:

$$\begin{aligned} u_d &= \frac{r_c^2 \gamma_w}{2k_h} \left(\frac{\partial \varepsilon_v}{\partial t} + \frac{k_v}{\gamma_w} \frac{\partial^2 \bar{u}_s}{\partial z^2} \right) \left\{ \frac{A^2}{A^2 - B^2} \left[\ln \frac{r}{r_c} - \frac{(A-B)F + (A+B)G}{2A} \right] + \frac{A}{2n^2 C^2} [(A+B)F + (A-B)G] \right\} + u_c(r_c), \\ u_n &= \frac{r_c^2 \gamma_w}{2k_h} \left(\frac{\partial \varepsilon_v}{\partial t} + \frac{k_v}{\gamma_w} \frac{\partial^2 \bar{u}_s}{\partial z^2} \right) \left\{ \left(\ln \frac{r}{r_s} - \frac{r^2 - r_s^2}{2r_c^2} \right) + \frac{A^2}{A^2 - B^2} \left[\ln s - \frac{1}{2} \left(\ln \alpha + \frac{BE}{A} \right) \right] + \frac{A^2}{2n^2 C^2} \left(\ln \alpha + \frac{BE}{A} \right) \right\} + u_c(r_c), \end{aligned} \quad (13)$$

式中, $E = \ln \frac{A+1}{A-1}$, $F(r) = \ln \frac{A+B - \frac{Cr}{r_c}}{A+1}$, $G(r) =$

$$\ln \frac{A-B + \frac{Cr}{r_c}}{A-1}.$$

再把式(13)代入式(5)可得:

$$\bar{u}_s = \frac{r_c^2 \gamma_w F_c}{2k_h} \left(\frac{\partial \varepsilon_v}{\partial t} + \frac{k_v}{\gamma_w} \frac{\partial^2 \bar{u}_s}{\partial z^2} \right) + u_c(r_c), \quad (14)$$

式中,

$$F_c = \frac{A^2 F_{c1} + n^2 F_{c2}}{n^2 - 1},$$

$$F_{c1} = \frac{1}{A^2 - B^2} \left(s^2 \ln s - \frac{s^2}{2} + \frac{1}{2} \right) - \frac{1}{(A^2 - B^2)C^2} \cdot$$

$$\left[-B + \frac{1}{2} - \left(\frac{A^2}{2} - B^2 \right) \ln \alpha + \frac{ABE}{2} \right] + \frac{1}{n^2 C^4} \left[-3B + \frac{1}{2} - \left(\frac{A^2}{2} + B^2 \right) \ln \alpha + \frac{3}{2} ABE \right],$$

$$F_{c2} = \ln \frac{n}{s} - \frac{3}{4} + \frac{s^2}{n^2} - \frac{s^4}{4n^4} + A^2 \left(1 - \frac{s^2}{n^2} \right) \cdot$$

$$\left[\frac{1}{A^2 - B^2} \left(\ln \frac{s}{\sqrt{\alpha}} - \frac{BE}{2A} \right) + \frac{1}{n^2 C^2} \left(\ln \sqrt{\alpha} - \frac{BE}{2A} \right) \right].$$

同理, 对式(10a)两边关于 r 积分并利用条件(11e)可得:

$$\frac{\partial u_c}{\partial r} = -\frac{\gamma_w r}{2k_{hc}} \left(\frac{\partial \varepsilon_v}{\partial t} + \frac{k_{vc}}{\gamma_w} \frac{\partial^2 \bar{u}_c}{\partial z^2} \right). \quad (15)$$

再对式(15)两边关于 r 积分可得:

$$u_c = u_c(r_c) + \frac{\gamma_w}{4k_{hc}} \left(\frac{\partial \varepsilon_v}{\partial t} + \frac{k_{vc}}{\gamma_w} \frac{\partial^2 \bar{u}_c}{\partial z^2} \right) (r_c^2 - r^2) \quad (16)$$

再把式(15)代入式(6)得:

$$\bar{u}_c = u_c(r_c) + \frac{\gamma_w r_c^2}{8k_{hc}} \left(\frac{\partial \varepsilon_v}{\partial t} + \frac{k_{vc}}{\gamma_w} \frac{\partial^2 \bar{u}_c}{\partial z^2} \right) \quad (17)$$

式(14)、(17)相减可得:

$$\bar{u}_s - \bar{u}_c = \left(\frac{r_c^2 F_c \gamma_w}{2k_h} - \frac{\gamma_w r_c^2}{8k_{hc}} \right) \frac{\partial \varepsilon_v}{\partial t} + \frac{r_c^2 F_c k_v}{2k_h} \frac{\partial^2 \bar{u}_s}{\partial z^2} - \frac{r_c^2 k_{vc}}{8k_{hc}} \frac{\partial^2 \bar{u}_c}{\partial z^2} \quad (18)$$

再把式(9)代入(18),消去 $\bar{u}_s$ 可得:

$$\bar{u} - \bar{u}_c = -\frac{\gamma_w (n^2 - 1)}{E_s (n^2 - 1 + Y)} \left(\frac{r_c^2 F_c}{2k_h} - \frac{r_c^2}{8k_{hc}} \right) \frac{\partial \bar{u}}{\partial t} + \frac{r_c^2 F_c k_v}{2k_h} \frac{\partial^2 \bar{u}}{\partial z^2} - \left[\frac{r_c^2 F_c k_v}{2k_h} + \frac{r_c^2 k_{vc} (n^2 - 1)}{8n^2 k_{hc}} \right] \frac{\partial^2 \bar{u}_c}{\partial z^2} \quad (19)$$

再把式(12a)、(15)代入式(11e)可得:

$$n^2 \frac{\partial \varepsilon_v}{\partial t} = - \left[(n^2 - 1) \frac{k_v}{\gamma_w} \frac{\partial^2 \bar{u}_s}{\partial z^2} + \frac{k_{vc}}{\gamma_w} \frac{\partial^2 \bar{u}_c}{\partial z^2} \right] \quad (20)$$

结合式(8)、(9)和(20)消去变量 $\bar{u}_s$ 有:

$$\frac{\partial^2 \bar{u}_c}{\partial z^2} = A' \frac{\partial \bar{u}}{\partial t} - B' \frac{\partial^2 \bar{u}}{\partial z^2} \quad (21)$$

式中, A' , B' 均为大于零的常系数,

$$A' = \frac{n^4 \gamma_w}{E_s (k_{vc} - k_v) (n^2 - 1 + Y)}, \quad B' = \frac{n^2 k_v}{k_{vc} - k_v}$$

再把式(21)代入式(19)整理得:

$$\bar{u}_c = \bar{u} + C' \frac{\partial \bar{u}}{\partial t} - D' \frac{\partial^2 \bar{u}}{\partial z^2} \quad (22)$$

式中, C' , D' 也均为大于零的常系数,

$$C' = \frac{\gamma_w [(n^2 - 1)k_{vc} + k_v]}{E_s (n^2 - 1 + Y) (k_{vc} - k_v)} \left[\frac{r_c^2 F_c}{2k_h} + \frac{r_c^2 (n^2 - 1)}{8k_{hc}} \right], \quad D' = \frac{k_v k_{vc}}{k_{vc} - k_v} \left[\frac{r_c^2 F_c}{2k_h} + \frac{r_c^2 (n^2 - 1)}{8k_{hc}} \right]$$

对式(22)关于 z 求二次导数,再把式(21)代入可得关于 $\bar{u}$ 的微分方程为

$$D' \frac{\partial^4 \bar{u}}{\partial z^4} - C' \frac{\partial^3 \bar{u}}{\partial t \partial z^2} - (B' + 1) \frac{\partial^2 \bar{u}}{\partial z^2} + A' \bar{u} = 0 \quad (23)$$

式(22)、(23)即为本文关于 $\bar{u}$ 和 $\bar{u}_c$ 的控制方程。

因为顶面排水,底面不排水,所以竖直方向边界条件可以写为

$$\left. \begin{aligned} \bar{u}(z, t) &= 0 \quad (z = 0), \\ \bar{u}_c(z, t) &= 0 \quad (z = 0), \end{aligned} \right\} \quad (24a)$$

$$\left. \begin{aligned} \frac{\partial \bar{u}(z, t)}{\partial z} &= 0 \quad (z = H), \\ \frac{\partial \bar{u}_c(z, t)}{\partial z} &= 0 \quad (z = H). \end{aligned} \right\} \quad (24b)$$

HAN等^[1]的研究发现:由于土体的排水性能远远低于桩体的排水性能,所以,在初始时刻,土体的不排水压缩模量远远大于桩体的排水压缩模量。正因如此,和砂井地基不同,初始时刻桩体内的孔压可能小于土体内的孔压。所以,传统的适用于砂井地基固结的初始条件已不能满足复合地基的固结问题。本文根据复合地基桩土共同承担外部荷载的特点,参照以前的研究^[6],认为在初始时刻,复合地基的平均初始孔压等于地基内的平均附加应力,即初始条件可以写为

$$\bar{u}(z, 0) = \sigma_0 \quad (25)$$

1.2 问题的解答

采用分离变量法,设方程(23)的解为 $\bar{u} = Z(z)T(t)$ 。

对方程(23)分离变量可得:

$$\frac{D'Z^{(4)} - (B' + 1)Z''}{A'Z - C'Z''} = -\frac{T'}{T} = \beta > 0 \quad (26)$$

式(26)可以写为

$$\left. \begin{aligned} D'Z^{(4)} - (B' + 1 - \beta C')Z'' - A'\beta Z &= 0, \\ T' + \beta T &= 0. \end{aligned} \right\} \quad (27)$$

参照TANG等^[10]的方法,可得式(21)的解的形式为

$$\bar{u} = \sum_{m=1}^{\infty} A_m [a_m \sin(\lambda_m z) + b_m \cos(\lambda_m z) + c_m \sinh(\xi_m z) + d_m \cosh(\xi_m z)] e^{-\beta_m t} \quad (28)$$

式中,

$$\left. \begin{aligned} \xi_m &= \sqrt{\frac{(B' + 1 - \beta_m C') + \sqrt{(B' + 1 - \beta_m C')^2 + 4A'D'\beta_m}}{2D'}}, \\ \lambda_m &= \sqrt{\frac{-(B' + 1 - \beta_m C') + \sqrt{(B' + 1 - \beta_m C')^2 + 4A'D'\beta_m}}{2D'}}. \end{aligned} \right\} \quad (29)$$

再把式(28)代入式(22)可得 $\bar{u}_c(z, t)$ 的表达式为

$$\bar{u}_c = \sum_{m=1}^{\infty} A_m \left\{ [1 - C'\beta_m + D'\lambda_m^2] [a_m \sin(\lambda_m z) + b_m \cos(\lambda_m z)] + (1 - C'\beta_m - D'\xi_m^2) [c_m \sinh(\xi_m z) + d_m \cosh(\xi_m z)] \right\} e^{-\beta_m t} \quad (30)$$

再由式(28)、(30)和边界条件(24)可得:

$$b_m = c_m = d_m = 0 \quad (31)$$

$$a_m \lambda_m \cos(\lambda_m H) = 0 \quad (32)$$

因为 $\lambda_m \neq 0$, a_m 也不能再为零(否则解为零解),所以只能有 $\cos(\lambda_m H)$ 为零,则可得:

$$\left. \begin{aligned} \lambda_m &= \frac{M}{H}, \\ M &= \frac{2m-1}{2} \pi \quad (m=1,2,3,\dots) \end{aligned} \right\} \quad (33)$$

再把式(33)代入式(29)得:

$$\beta_m = \frac{c_v(n^2-1+Y)}{r_c^2} \left\{ \frac{n^4}{M^2} \left(\frac{H}{r_c} \right)^2 + \left[(n^2-1) \frac{k_{vc}}{k_v} + 1 \right] \cdot \left[\frac{n^2 F_c k_{vc}}{2 k_h} + \frac{(n^2-1) k_{vc}}{8 k_{hc}} \right] \right\}^{-1} \left\{ \left[\frac{n^2 F_c k_v}{2 k_h} + \frac{(n^2-1) k_v}{8 k_{hc}} \right] \cdot \left(\frac{M r_c}{H} \right)^2 + \left[(n^2-1) + \frac{k_{vc}}{k_v} \right] \right\}. \quad (34)$$

把式(34)代回式(28)、(30)并利用初始条件方程最终的解答为

$$\bar{u} = \sigma_0 \sum_{m=1}^{\infty} \frac{2}{M} \sin \left(\frac{M}{H} z \right) e^{-\beta_m t}, \quad (35)$$

$$\bar{u}_c = \sigma_0 \sum_{m=1}^{\infty} \frac{2}{M} \left(1 - C' \beta_m + D' \frac{M^2}{H^2} \right) \sin \left(\frac{M}{H} z \right) e^{-\beta_m t}. \quad (36)$$

2 固结度

2.1 固结度的表达式

一般来说, 复合地基的固结度可以根据应力来定义, 也可以根据应变来定义。参考文献[6]的方法可以证明, 本文按应力和按应变定义的固结度相等, 所以, 此处只给出按应力定义的固结度的表达式。

因为本文考虑了桩土共同作用以及桩体的固结压缩, 所以, 在计算复合地基的固结度时, 把桩体的固结压缩也引入地基的固结度计算中, 即地基任一时刻按应力定义的固结度可以表示为复合地基任一时刻的总平均有效应力和总应力之比, 即:

$$U(t) = \frac{\int_0^H (\sigma_0 - \bar{u}) dz}{\int_0^H \sigma_0 dz} = 1 - \frac{\int_0^H \bar{u} dz}{\int_0^H \sigma_0 dz}, \quad (37)$$

把式(35)代入式(37)得:

$$U(t) = 1 - \sum_{m=1}^{\infty} \frac{2}{M^2} e^{-\beta_m t}. \quad (38)$$

2.2 讨论

当 $n \rightarrow 1$ 时, 固结度的表达式(38)退化为

$$U(t) = 1 - \sum_{m=1}^{\infty} \frac{2}{M^2} e^{-M^2 T_{vc}}, \quad (39)$$

式中, T_{vc} , c_{vc} 分别为桩的竖向时间因子和竖向固结系数, $T_{vc} = \frac{c_{vc} t}{H^2}$, $c_{vc} = \frac{k_{vc} E_c}{\gamma_w}$ 。

显然, 式(39)为 Terzaghi 的天然地基一维固结

解。实际上, 当 $n \rightarrow 1$ 时, 土体整个被桩体置换, 这时候复合地基即转化为天然地基。另外, 当 $n \rightarrow \infty$ 时, $F_c \rightarrow \infty$, 这时固结度的表达式(38)退化为

$$U(t) = 1 - \sum_{m=1}^{\infty} \frac{2}{M^2} e^{-M^2 T_v}, \quad (40)$$

式中, T_v , c_v 分别为土体的竖向时间因子和竖向固结系数, $T_v = \frac{c_v t}{H^2}$, $c_v = \frac{k_v E_s}{\gamma_w}$ 。

同样, 式(40)也是 Terzaghi 的天然地基一维固结解。不同的是, 其时间因子为土体的竖向时间因子。实际上, 当 $n \rightarrow \infty$ 时, 相当于桩体的半径趋于 0, 此时, 复合地基也必然转化为天然地基。

3 解的比较

由图3可以看出, 随着井径比的增大, 本文解和文献[10]的解差值减小。实际上, 井径比的增大意味着桩体直径减小, 这时考虑桩体内径向渗流和不考虑桩体内径向渗流对地基的固结性状的影响也就变小。

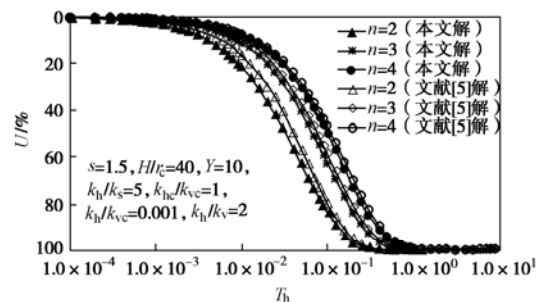


图3 不同井径比时解的比较

Fig. 3 Comparison among solutions under different diameter ratios

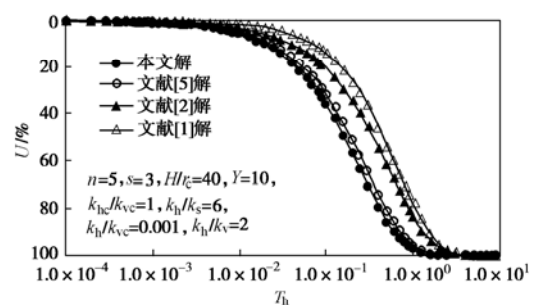


图4 和以往解的比较

Fig. 4 Comparison between present and previous solutions

由图4可以看出, 本文考虑扰动区土体水平渗透系数抛物线变化的解固结最快, 文献[5]考虑扰动区土体水平渗透系数直线变化的解固结次之, 文献[1, 2]考虑扰动区水平渗透系数不变的解最慢, 而文献[1]忽略了土体里的竖向渗流, 其解又慢于文献[2]的解。需要说明的是, 图4中本文解和文献[5]解的差异主要是扰动区土体水平渗透系数不同的结果, 考虑或者不考

虑桩体内径向渗流在井径比较大时对固结度影响较小(见图3)。

4 结 论

(1) 不管井径比趋于 1 或无穷大, 本文的解均能退化到 Terzaghi 天然地基一维固结解。

(2) 考虑扰动区土体水平渗透系数抛物线变化比考虑直线变化或不变所得的地基固结速率要快。

(3) 随着井径比的减小, 考虑桩体内径向渗流和不考虑桩体内径向渗流对地基的固结性状影响变大。

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